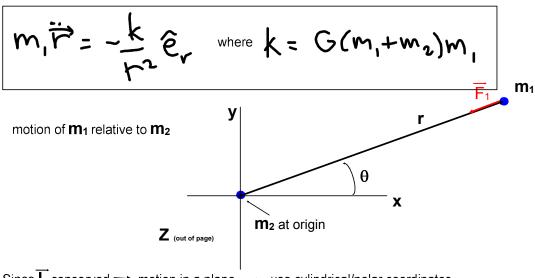
Solution to motion under a central force (inverse square law):



Since \overline{L} conserved \longrightarrow motion in a plane \longrightarrow use cylindrical/polar coordinates

unit vectors:
$$\hat{e}_{r}$$
 \hat{e}_{a} \hat{e}_{z}

Write the D.E. (top of page) in cyl/polar coordinates:

Note, from vector calculus:

The vector equation at the top of the page can now be written as 2 scalar DE's, one for each polar component:

$$\hat{C}_{\bullet}: \quad m_{\bullet}(\dot{r} - r\dot{\theta}^{2}) = -\frac{k}{r^{2}}$$

$$\hat{C}_{\bullet}: \quad m_{\bullet}(2\dot{r}\dot{\theta} + r\dot{\theta}) = 0$$

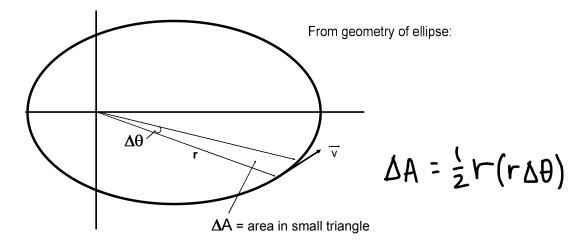
Solve to get r = r(t) $\theta = \theta(t)$ then eliminate \mathbf{t} to get $\mathbf{r} = \mathbf{r}(\boldsymbol{\theta})$ for shape of orbit:

One possible solution is an elliptical orbit with one focus at the origin (Law of Ellipses, Kepler's 1st law)

Solve the " θ " equation for Kepler's 2nd law (Law of Areas):

$$\begin{aligned} & \text{M}_{1}\left(2r\dot{\theta}+r\dot{\theta}\right)=0 & \text{Note:} & \frac{d}{dt}\left(r^{2}\dot{\theta}\right)=2r\dot{\theta}r+r^{2}\dot{\theta} \\ & = r\left(2\dot{\theta}\dot{r}+r\dot{\theta}\right) \end{aligned}$$

$$\text{Then:} & \text{M}_{1}\frac{d}{dt}\left(r^{2}\dot{\theta}\right)=0 & \frac{d}{dt}\left(r^{2}\dot{\theta}\right)=0 \\ & \text{Then:} & \text{Mote:} & \frac{d}{dt}\left(r^{2}\dot{\theta}\right)=0 \\ & \text{Mote:}$$



$$\lim_{A\to 0} \frac{\Delta A}{\Delta t} = \frac{1}{2} r^2 \frac{\Delta \theta}{\Delta t} \longrightarrow \frac{dA}{dt} = \frac{1}{2} r^2 \dot{\theta} = \text{"areal velocity"}$$

$$r^2 \dot{\theta} = 1 \text{ and } 1 = \text{const} \longrightarrow \frac{dA}{dt} = \frac{1}{2} l = \text{const}$$

Equal areas in equal time intervals: Law of Areas

General solution (after eliminating **t**):

Compare to the polar equation for conic sections:

